

CS 101

Numerical Geometric Integration

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Today: More Symplecticity

We've seen some examples

- Symplectic Euler, Stormer-Verlet, implicit midpoint

Generating Functions

- Continuous & discrete picture

Variational Integrators

- Another approach to symplecticity

Condition of Symplecticity

$$(p, q) \rightarrow (\hat{P}(p, q), \hat{Q}(p, q))$$

Initial condition Final state

When is it symplectic?

$$\left(\frac{\partial(\hat{P}, \hat{Q})^T}{\partial(p, q)} \right) J \left(\frac{\partial(\hat{P}, \hat{Q})}{\partial(p, q)} \right) = J$$

Let's (i.e. you) write it out...

Condition of Symplecticity

$$\hat{P}_p^T \hat{Q}_p = \hat{Q}_p^T \hat{P}_p$$

$$\hat{Q}_q^T \hat{P}_q = \hat{P}_q^T \hat{Q}_q$$

$$\hat{P}_p^T \hat{Q}_q - I = \hat{Q}_p^T \hat{P}_q$$

Just says these quantities are symmetric

$(\hat{P}(p, q), \hat{Q}(p, q))$ is symplectic iff these are true

Condition of Symmetry

Out of the blue:

$$f(p, q) = \hat{P}^T \nabla \hat{Q} - p^T \nabla q$$

If $\nabla f(p, q)$ is symmetric then*

$$f(p, q) = \nabla S(p, q)$$

for some S

- When is this $\nabla f(p, q)$ symmetric?

➤ First step: $\nabla \hat{Q} = \hat{Q}_p \nabla p + \hat{Q}_q \nabla q$

Same Conditions

Symmetry of $\nabla f(p, q)$ same as symplecticity of (\hat{P}, \hat{Q}) !

So what?

- $S(p, q)$ exists iff map is symplectic

Change of variable*: $S(q, \hat{Q})$

$$\nabla S(q, \hat{Q}) = \frac{\partial S}{\partial q} \nabla q + \frac{\partial S}{\partial \hat{Q}} \nabla \hat{Q}$$

Generating Function

$$\begin{aligned}\nabla S(q, \hat{Q}) &= \frac{\partial S}{\partial q} \nabla q + \frac{\partial S}{\partial \hat{Q}} \nabla \hat{Q} \\ &= \hat{p}^T \nabla \hat{Q} - p^T \nabla q = f(p, q) \\ \hat{p}^T &= \frac{\partial S}{\partial \hat{Q}}(q, \hat{Q}) \quad p^T = -\frac{\partial S}{\partial q}(q, \hat{Q})\end{aligned}$$

- Any [smooth] $S(q, \hat{Q})$ "generates" a symplectic map $(p, q) \rightarrow (\hat{p}(p, q), \hat{Q}(p, q))$
- Similar for $S(q, \hat{p}), S(p, \hat{Q}),$ etc.

How is this useful?

One useful variant: $S(q, \hat{p})$

$$\hat{Q}^T = \frac{\partial S}{\partial \hat{p}}(q, \hat{p}) \quad p^T = \frac{\partial S}{\partial q}(q, \hat{p})$$

Take: $S = \bar{S} + \hat{p}^T q$

$$\hat{Q}^T = q^T + \frac{\partial \bar{S}}{\partial \hat{p}}(q, \hat{p}) \quad \hat{p}^T = p^T - \frac{\partial \bar{S}}{\partial q}(q, \hat{p})$$

Try $\bar{S}(q, \hat{p}) = hH(q, \hat{p})$ We get symplectic Euler

Creating More Integrators

Choice of S gives integrator

- Symplectic euler, midpoint, RK
- More clever choices
 - See book
- All symplectic integrators have a generating function.

That Was Not Easy

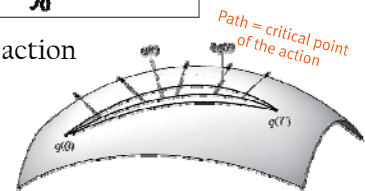
We're going to see this again...

Reminder

Physical path extremizes action

$$\delta S(q) = \delta \int_0^T L(q, \dot{q}) dt = 0$$

- stationary action

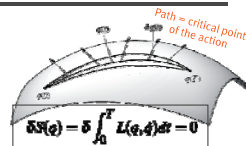


Variational Mechanics

Hamilton's Principle

Lagrangian: $L(q, \dot{q}) = K(\dot{q}) - W(q)$

Action: $S(q) = \int_0^T L(q, \dot{q}) dt$



$$\begin{aligned}\delta S(q) &= \delta \int_0^T L(q(t), \dot{q}(t)) dt = \int_0^T \left[\frac{\partial L}{\partial q} \cdot \delta q + \frac{\partial L}{\partial \dot{q}} \cdot \delta \dot{q} \right] dt \\ &= \int_0^T \left[\frac{\partial L}{\partial q} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) \right] \delta q dt + \left[\frac{\partial L}{\partial \dot{q}} \delta q \right]_0^T,\end{aligned}$$

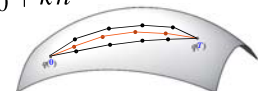
Euler-Lagrange Equation \Rightarrow Newton's Second Law

Discrete Variational Mechanics

Hamilton's Stationary Action Principle:

a discrete dynamical system always finds an optimal course from one position to another

- discrete time $t_k = t_0 + kh$
- so path = "polyline"
- use same recipe!
 - but with variations of states $q_k := q(t_k)$



Discrete Variational Mechanics

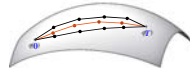
Hamilton's Stationary Action Principle:

- Discrete Lagrangian

$$\int_{t_k}^{t_{k+1}} L(q, \dot{q}) dt \approx L_d(q_k, q_{k+1})$$

- obtained, for instance, through quadrature

$$\triangleright L_d(q_k, q_{k+1}) = hL\left(\frac{q_k + q_{k+1}}{2}, \frac{q_{k+1} - q_k}{h}\right)$$



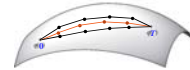
Discrete Variational Mechanics

Hamilton's Stationary Action Principle:

- Discrete Action

$$S_d(q) = \sum_{k=0}^N L_d(q_k, q_{k+1})$$

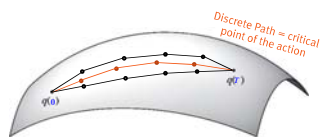
- integral of L, like before $S(q) = \int_0^T L(q, \dot{q}) dt$



Discrete Variational Mechanics

Hamilton's Stationary Action Principle:

$$\delta S_d(q) = \delta \sum_{k=0}^N L_d(q_k, q_{k+1}) = 0$$



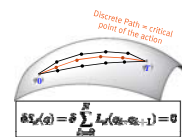
Discrete Variational Mechanics

Turning the crank:

$$D_1 L_d(q_{k+1}, q_{k+2}) + D_2 L_d(q_k, q_{k+1}) = 0$$

Discrete Euler-Lagrange

derivatives with respect to first or second variable



- must be true for each q_{k+1}
- so in particular...

$$\text{solve DEL for time integration: } (q_k, q_{k+1}) \rightarrow q_{k+2}$$

Summary

So Time Integration's a Breeze

- start from q_0 and q_1
 - can start with momentum instead too
- use DEL $(q_k, q_{k+1}) \rightarrow q_{k+2}$
- iterate...

Discrete path optimal btw 0 and T

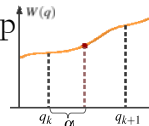
- discrete Hamilton's principle

Discrete Lagrangian: More Details

$$\int_{t_k}^{t_{k+1}} L(q, \dot{q}) dt = \int_{t_k}^{t_{k+1}} [K(\dot{q}) - W(q)] dt$$

$$L_d(q_k, q_{k+1}) = h \left[\frac{1}{2} \left(\frac{q_{k+1} - q_k}{h} \right)^T M \left(\frac{q_{k+1} - q_k}{h} \right) - W((1 - \alpha)q_k + \alpha q_{k+1}) \right]$$

- integral over one time step
- use any quadrature
 - trapezoidal, midpoint, ...
- discretize potential W



Simple Discrete Example

Let's try the free-falling mass again:

$$L_d(q_k, q_{k+1}) = h \left(\frac{1}{2} m \frac{(q_{k+1} - q_k)^2}{h} - mg \frac{q_k + q_{k+1}}{2} \right)$$

$$D_2 L_d(q_k, q_{k+1}) = m \frac{q_{k+1} - q_k}{h} - \frac{hmg}{2}$$

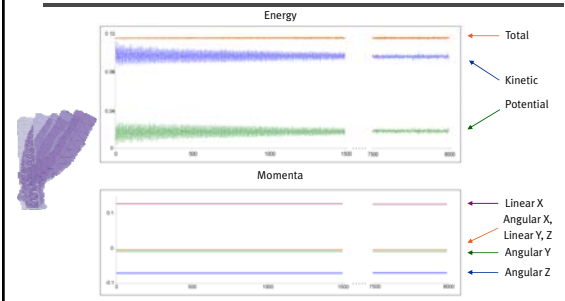
$$D_1 L_d(q_{k+1}, q_{k+2}) = -m \frac{q_{k+2} - q_{k+1}}{h} - \frac{hmg}{2}$$

So, we get:

$$D_1 L_d(q_{k+1}, q_{k+2}) + D_2 L_d(q_k, q_{k+1}) = 0 \Leftrightarrow \frac{q_{k+2} - 2q_{k+1} + q_k}{h^2} = g$$

$$\ddot{q} = g$$

Conservation Laws



Where's Waldo?

The momentum p disappeared

- or did it?

Legendre transform to the rescue

- see slide 11 of lecture 3: $p = \frac{\partial \mathcal{L}(q, \dot{q})}{\partial \dot{q}}$
- here: $p_n = -D_1 L(q_n, q_{n+1})$
- allows you to switch: $(p_n, q_n) \rightarrow (q_n, q_{n+1})$
- perfect for initial conditions...

DEL in Momentum/Position

We can rewrite the DEL equations as:

$$p_n = -D_1 L(q_n, q_{n+1}) = D_2 L(q_{n-1}, q_n)$$

- rewrite one more time

$$p_{n+1} = \frac{\partial L}{\partial q_{n+1}}(q_n, q_{n+1}) \quad p_n = -\frac{\partial L}{\partial q_n}(q_n, q_{n+1})$$

$$\hat{p}^T = \frac{\partial S}{\partial \hat{Q}}(q, \hat{Q}) \quad p^T = -\frac{\partial S}{\partial q}(q, \hat{Q})$$

L is a generating function so these are symplectic!